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The Stability Of Crank – Nicholson and Explicit Methods for Numerical **Solution for Sine – Gordon Equation**

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ABSTRACT

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In This Paper, We Solve The Sine – Gordon Equation by two Numerical Methods: Crank - Nicholson and Explicit and we discuss The Stabilities, and we obtained That The Stability Crank - Nicholson Methods is More Than The Stability Of Explicit Methods.

استقراريه الحل العددي بطريقتي كرانك -نيكلسون Crank-Nicholson والطريقة الصريحة Explicit Method لمعادلة Sine-Gordon

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الملخص

في هذا البحث قمنا بحل معادلة Sine-Gordon باستخدام طريقتين عديبتين .هما طريقة كرانك – نيكلسون . (Crank-Nicholson) وطريقة Explicit Method وناقشنا استقرارية الحلول لكل منهما ووجدنا أن طريقة كرانك – نيكلسون أكثر استقراراً من الطريقة الصريحة. الكلمات المفتاحية: الاستقرارية، معادلة Sine-Gordon، طربقة Crank-Nicholson، الطريقة الصريحة (Explicit Method)

Introduction

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The differential equation is one of the types to describe many of Phenomena in Physics, Chemical , Engineer . and there are Many numerical Methods that are Used to Solving the Problems, Finite difference Methods (FDM) and Finite Elements Methods (FEM)

are two Methods are More using , but . There are a lot of risks When applying These Methods [1]. The numerical Solution for Partial differential equations is absultly find The experimental solution specially When We Using (FDM), [2].

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In (1939) Frenkel, Kontorova are in traduced Sine – Gordon equation in Physics. And in (1962) Peering and Skyrme are Studied the numerical result for some Chemical Problems [11]. Forest and Mclaughlin, Ercolani in (1990) are improved the geometry formula of Sine – Gordon equation which can be integrable under periodically boundary Conditions, and they are obtain the good numerical results [5].

2- The Mathematical Model

Now, we derive the mathematical Model for Sine-Gordon equation by Klein – Gordon equation, where the governing equation is in the bellow formula

$$\frac{\partial^2 u}{\partial t^2} - d^2 \,\frac{\partial^2 u}{\partial x^2} + f(u) = 0 \tag{1}$$

Where d is the wave speed, f(u) the density of elasticity power.

Note: if f(u) = g(u) then, the equation (1) becomes Klein – Gordon linear equation

i.e.
$$:\frac{\partial^2 u}{\partial t^2} - d^2 \frac{\partial^2 u}{\partial x^2} + gu = 0$$
 (2)

if g = 0 then Klein – Gordon become traditional Wave equation , and if $f(u) = \sin u$ the equation (2) become Sine – Gordon equation [5], [8].

$$\frac{\partial^2 u}{\partial t^2} - d^2 \frac{\partial^2 u}{\partial x^2} = -\sin u \tag{3}$$

$$\frac{\partial^2 u}{\partial t^2} - \frac{\partial^2 u}{\partial x^2} = -\sin u \tag{4}$$

3- Deriving the general form of Explicit Method

In (2003) Ares, Guesta, Sanchez and Toral are study the model of Sine – Gordon equation analytically and numerically[4].

In (2005)Khomeriki and Leon are described the Bi – Stability behavior Case as appositive non linear equation and they are explain this numerically [7].

And with boundary and initial Conditions

$$u(x, 0) = p + e_0 \cos[mx] , u_t(x, 0) = 0$$
$$m \frac{1}{\sqrt{2}} , l = 2\sqrt{2}p , 0 \le e_o \le 1000 ,$$
$$-L \le x \le L$$

u(-L,t) = u(L,t) = b, where b is constant

The following equation is called [Perturbed Sine Gordon Equation]

$$\frac{\partial^2 u}{\partial t^2} - \frac{\partial^2 u}{\partial x^2} = -\sin u + e \sin[d_t + d_o] \quad (5)$$

Where e is a capacity and d is frequency, d_o is a Phase [10].

Also, the equation bellow is called Fractional Sine – Gordon equation

$$u_{xx} - {}^R D^a u + Sin u = 0 \tag{6}$$

 $^{R} D^{a}$ is the Fractional derivative on Riesz space

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First Step we divide the region. Rectangles: $R = [(x, t) = -L \le x \le L, 0 \le t \le C]$ to [n-1], [m-1] from the Rectangles for Length Side $\Delta x = h$, $\Delta t = k$, as show bellow [9]



The form (1) represent mesh (Explicit method)

From level 1: Calculate the value $t = t_1 = 0$ [3]

 $u(x_i, t_1) = p + e_o \cos[mx_i]$, i = 1, 2, 3, 4, ..., n - 1 where

$$0 \leq e_o \leq 1000$$
 , $m = \frac{1}{\sqrt{2}}$, $l = 2 \sqrt{2} \, p$

Where the differences are used to approximate u_{tt} (x,t), $u_{xx}(x,t)$

$$u_{tt}(x,t) = \frac{u(x,t+k) - 2u(x,t) + u(x,t-k)}{k^2} + O(k^2)$$
(7)

$$u_{xx}(x,t) = \frac{u(x+h,t)-2u(x,t)+u(x-h,t)}{h^2} + 0 (h^2)$$
(8)

And the nodes are :

 $x_{i+1} = x_i + h$, $x_{i-1} = x_i - h$, $t_{j+1} = t_j + k$, $t_{j-1} = t_j - k$

Ignore $O(k^2)$, $O(h^2)$ the we obtain getting:

$$u_{tt}(x,t) = \frac{u(x,t+k) - 2u(x,t) + u(x,t-k)}{k^2}$$
(9)

$$u_{xx}(x,t) = \frac{u(x+h,t)-2u(x,t)+u(x-h,t)}{h^2}$$
(10)

Also , to find the values 2^{nd} ($t = t_2$) level by Taylor series [9]:

$$u(x,t) = u(x,0) + u_t(x,0)k + \frac{u_{tt}(x,0)k^2}{2!} + O[k^3]$$
(11)

Put
$$x = x_i$$

$$\frac{u(x_{i},t) = u(x_{i},0) + u_{t}(x_{i},0)k + u_{t}(x_{i},0)k + u_{t}(x_{i},0)k + u_{t}(x_{i},0)k^{2}}{2} + O[k^{3}]$$
(12)

Where

$$u_{tt}(x_{i},0) = \frac{u(x_{i+1},0) - 2u(x_{i},0) + u(x_{i-1},0)}{h^{2}} - \sin[u(x_{i},0)]$$
(13)

And substitute equations (13,14) in equation (12) we get :

$$u(x_{i},k) =$$

$$u(x_{i},0) + \left[\frac{u(x_{i+1},0) - 2u(x_{i},0) + u(x_{i-1},0)}{h^{2}} - sin[u(x_{i},0)]\right]\frac{k^{2}}{2} + O(k^{3})$$

$$u(x_{i},k) = u(x_{i},0) + \frac{r^{2}}{2}[u(x_{i+1},0) - sin(x_{i+1},0)] - sin(x_{i},k) = u(x_{i},0) + \frac{r^{2}}{2}[u(x_{i+1},0) - sin(x_{i},0)]$$

$$u(x_{i}, \kappa) = u(x_{i}, 0) + \frac{1}{2} [u(x_{i+1}, 0) - \frac{1}{2} [u(x_{i+1}, 0)] - \frac{1}{2} [u(x_{i}, 0)] - \frac{k^{2}}{2} sin[u(x_{i}, 0)]$$
(15)

And with $r = \frac{k}{h}$, and substitute ($u_{i,1}$) of with $u(x_i, 0)$, and with taking operation simplied to equation (15) we get on level two (FDM)

$$u_{i,2} = [1 - r^2]u_{i,1} + \frac{r^2}{2} [u_{i+1,1} + u_{i-1,1}] - \frac{k^2}{2} \sin[u_{i,1}]$$
(16)

Now show that Calculate value u(x, t) as nodes points for class other :

$$u(x_i, t_j): i = 1, 2, 3, ..., n-1$$
, $j = 2, 3, 4, ..., m$

As substitute $(u_{i,j})$ with out $u(x_i, t_j)$ in equations (10, 11), and by substitute equations complete in equation (3), we get :

$$\frac{u_{i,j+1}-2 u_{i,j}+u_{i,j-1}}{k^2} - \frac{u_{i+1,j}-2 u_{i,j}+u_{i-1,j}}{h^2} = \sin[u_{i,j}]$$
(17)

And from equation (17) we get :

$$u_{i,j+1} = -u_{i,j-1} + [2 - 2r^2] u_{i,j} + r^2 [u_{i+1,j} + u_{i-1,j}] - k^2 \sin[u_{i,j}]$$
(18)

And equation (18) represent (FDM) by used Explicit Method for equation Sine - Gordon and used equation (18) Calculate (j + 1), and by value (j - 1, j), these the method Calculate form Explicit for value undefined $(u_{i,j+1})$, and by value defined

$$u_{i,j-1}$$
 , $u_{i+1,j}$, $u_{i,j}$, $u_{i-1,j}$

That show in form (1)

4- The Stability of Explicit Method

Now . we explain the Stability of Explicit Method with Von- Neumann the solution step are

$$u_{x,t}=\,e^{gt}\,.\,e^{ibx}$$
 , $x=nh$, $t=mk$, $h=\Delta x$, $k=\,\Delta t$

TIPS

$$g\,,b\,>0$$
 , $i=\sqrt{-1}$

 $\begin{array}{l} u_{x,t} = e^{gmk} \cdot e^{ibnh} = \ [e^{g\Delta t}]^m \ e^{ibn\Delta x} = \\ x^m \ e^{ibn\Delta x} \end{array}$

The main step that is substitute the solution of (FDM) at time t with $x^m e^{ibnDx}$ [12] and with Von-Neumann can be used [Linearized Stability Analysis] [6]:

$$\frac{\partial^2 u}{\partial t^2} = \frac{\partial^2 u}{\partial x^2} - u \tag{19}$$

And by explicit Method

$$u_{n,m+1} = -u_{n,m-1} + [2 - 2r^2]u_{n,m} + r^2[u_{n+1,m} + u_{n-1,m}] - [\Delta t]^2 u_{n,m}$$
(20)

And by substitute $u_{n,m} = x^m e^{ibn\Delta x}$ in equation (20) we geting

 $x^{m+1} \cdot e^{ibn\Delta x} = -x^{m-1} e^{ibn\Delta x} + [2 - 2r^2]x^m e^{ibn\Delta x} + r^2 [x^m e^{ib[n+1]\Delta x} + x^m e^{ib[n-1]\Delta x}] - [\Delta t]^2 x^m e^{ibn\Delta x}$

$$\begin{split} x^m x . e^{ibn\Delta x} &= -x^m x^{-1} e^{ibn\Delta x} + [2 - 2r^2 - (\Delta t)^2] x^m e^{ibn\Delta x} + r^2 [x^m e^{ib n\Delta x} e^{ib\Delta x} + x^m e^{ib n\Delta x} e^{-ib\Delta x}] \end{split}$$

And by divide on $x^m e^{ibnDx}$ we getting

$$\begin{aligned} x &= -x^{-1} + \left[2 - 2r^2 - (\Delta t)^2\right] + r^2 \left[e^{ib\,\Delta x} + e^{-ib\,\Delta x}\right] \rightarrow \\ \frac{x^{2}+1}{x} &= \left[2 - 2r^2 - (\Delta t)^2\right] + 2r^2 \cos\left[b\,\Delta x\right] \rightarrow \\ \frac{x^{2}+1}{x} &= \left[2 - 2r^2 - (\Delta t)^2\right] + 2r^2 \left[1 - 2Sin^2\left(\frac{b\,\Delta x}{2}\right)\right] \rightarrow \\ \frac{x^{2}+1}{x} &= 2\left[1 - \frac{(\Delta t)^2}{2} - 2r^2 Sin^2\left(\frac{b\,\Delta x}{2}\right)\right] \rightarrow \\ \frac{x^{2}+1}{x} &= 2a \\ a &= 1 - \frac{(\Delta t)^2}{2} - 2r^2 Sin^2\left(\frac{b\,\Delta x}{2}\right) \\ x^2 - 2ax + 1 &= 0 \implies x = am\sqrt{a^2 - 1} \\ \text{The must } |a| \leq 1 \text{ , x is factor} \end{aligned}$$

$$\left|1 - \frac{(\Delta t)^2}{2} - 2r^2 \operatorname{Sin}^2\left(\frac{b\Delta x}{2}\right)\right| \le 1 \quad (21)$$

And from Inequalities (21), we getting

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<u>https://doi.org/10.25130/tjps.v28i4.1532</u> −1 ≤ 1 − $\frac{(\Delta t)^2}{2}$ − $2r^2 Sin^2 \left(\frac{b\Delta x}{2}\right) \le 1$

And by taking the side Inequality (21), we getting right

$$1 - \frac{(\Delta t)^2}{2} - 2r^2 \operatorname{Sin}^2\left(\frac{b\Delta x}{2}\right) \le 1$$
 (22)

And from Inequalities (22), we getting

$$r^2 \geq \frac{-(\Delta t)^2}{4 \operatorname{Sin}^2(\frac{b \,\Delta x}{2})}$$

Since $r^2 = \frac{(Dt)^2}{(Dx)^2}$, the value r^2 is positive, Inequality (22) way to $r^2 > 0$ this true early, for to suites Inequality (21), we need to:

$$-1 \le 1 - \frac{(\Delta t)^2}{2} - 2r^2 \sin^2\left(\frac{b\Delta x}{2}\right)$$
$$2 \ge \frac{(\Delta t)^2}{2} + 2r^2 \sin^2\left(\frac{b\Delta x}{2}\right)$$

$$1 - \frac{\left(\Delta t\right)^2}{2} \ge r^2 \operatorname{Sin}^2\left(\frac{b\Delta x}{2}\right)$$

And since $Sin^2\left(\frac{b\Delta x}{2}\right)$ equal 1 for value b :

$$r^2 \leq \frac{4 - (\Delta t)^2}{4}$$
 (23)

And since $r^2 = \frac{(\Delta t)^2}{(\Delta x)^2}$ And from Inequality (22), we getting

$$(\Delta t)^2 \leq \frac{4 \ (\Delta x)^2}{4 + (\Delta x)^2} \tag{24}$$

And represent Inequalities (23,24) on Explicit Method , must the Stability no operation by Sine-Gordon equation .

5- Deriving The General formula of [Crank – Nicholson]

In [*Crank* – *Nicholson*] method we replace the 2^{nd} partial derivative u_{xx} by the average of central difference at time t and points (j - 1, j + 1) as shown in the figure (2) bellow [13] :



The form (2) show that solution by using Crank - Nicholson

$$\frac{\frac{\partial^2 u}{\partial x^2}}{\frac{1}{2} \left(\frac{u_{i+1,j+1} - 2u_{i,j+1} + u_{i-1,j+1}}{h^2} + \frac{u_{i+1,j-1} - 2u_{i,j-1} + u_{i-1,j-1}}{h^2}\right)}{(25)$$

$$\frac{\partial^2 u}{\partial x^2} = \frac{u_{i,j+1} - 2u_{i,j} + u_{i,j-1}}{k^2}$$
(26)

And as substitute the equation (25,26) in equation (3) we get :

$$\frac{u_{i,j+1}-2u_{i,j}+u_{i,j-1}}{k^2} + \frac{u_{i+1,j+1}-2u_{i,j+1}+u_{i-1,j+1}}{2h^2} - \frac{u_{i+1,j-1}-2u_{i,j-1}+u_{i-1,j-1}}{2h^2} = \sin[u_{i,j}]$$
(27)

And from equation (28) we get :

$$-r^{2} \left[u_{i+1,j+1} + u_{i-1,j+1} \right] + \left[2 + 2r^{2} \right] u_{i,j+1}$$
$$= 4 u_{i,i} - 2k^{2} \sin[u_{i,i}] + r^{2} \left[u_{i+1,i-1} + \right]$$

$$= 4 u_{i,j} - 2k \sin[u_{i,j}] + 7 [u_{i+1,j-1} + u_{i-1,j-1}] - [2 + 2r^2] u_{i,j-1}$$

$$(28)$$

Where the equation (28) is the approximate the finite difference of [Crank - Nicholson] method for [Sine - Gordon] method, and hence we have : three variables $u_{i+1,j+1}$, $u_{i,j+1}$, $u_{i-1,j+1}$ and

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all terms in right hand side of equation (28) as know, so the equation (28) can be lead to the algebraic liner system with three diameters

$$AX = B \tag{29}$$

A: is a matrix have three diameters , X: Vertical vector , B: Vertical vector

Example: Let we have rectangular plate (2m, 4m) and it is coincide on x and y axis at original point

$$\Delta t = 4 \ sec$$
, $\Delta x = 5 \ cm$, $a = 1.25$

Initial condition $v(x, t) = 0^{\circ} \quad 0 < x < 10 \quad t = 0$

Boundary condition $v(0,t) = 8^\circ$, $v(10,t) = 4^\circ$

Solution:

1) Make the mesh and substitute the values on shapes



2) The crank – Nicolson general formula

$$-rv_{l-1}^{i+1} + 2(1+m)v_l^{i+1} - rv_{l+1}^{i+1}$$
$$= rv_{l-1}^i + 2(1-m)v_l^i + rv_{l+1}^i$$

To find m

$$\begin{split} m &= a \frac{\Delta t}{\Delta x^2} \rightarrow r = 1.25 \frac{5}{5^2} = 0.2 \\ \text{Substitute from l} = 123, i = 0, m = 0.2 \\ &-0.2v_0^1 + 2.4v_1^1 - 0.2v_2^1 \\ &= 0.2v_0^0 + 1.6v_1^0 + 0.2v_2^0 \\ &-0.2v_1^1 + 2.4v_2^1 - 0.2v_3^1 \\ &= 0.2v_1^0 + 1.6v_2^0 + 0.2v_3^0 \\ &-0.2v_2^1 + 2.4v_3^1 - 0.2v_4^1 \\ &= 0.2v_2^0 + 1.6v_3^0 + 0.2v_4^0 \\ \\ \text{Substitute initial and boundary condition, we get} \\ &-0.2(8) + 2.4v_1^1 - 0.2v_2^1 \\ &= 0.2(0) + 1.6(0) + 0.2(0) \\ &-0.2v_1^1 + 2.4v_2^1 - 0.2v_3^1 \\ &= 0.2(0) + 1.6(0) + 0.2(0) \\ &-0.2v_2^1 + 2.4v_3^1 - 0.2(4) \\ &= 0.2(0) + 1.6(0) + 0.2(0) \end{split}$$

Hence, we have $-0.2v_1^1 + 2.4v_2^1 - 0.2v_3^1 = 0 \dots \dots \dots (2)$ $-0.2v_2^1 + 2.4v_3^1 = 0.8 \qquad \dots \dots \dots (3)$ Multiply equation (2) by 12 and solve it with equation (1) we have $-2.4v_1^1 + 28.8v_2^1 - 2.4v_3^1 = 0$ $2.4v_1^1 - 0.2v_2^1 = 1.6$ $28.6v_2^1 - 2.4v_3^1 = 1.6 \qquad \dots \dots \dots (4)$ Solve equation (4) and adding to equation (3) $28.6v_2^1 - 2.4v_3^1 = 1.6$ $-0.2v_2^1 + 2.4v_3^1 = 0.8$ $28.4v_2^1 = 2.4$ $v_2^1 = 0.0845070423$ Substitute v_2^1 value in equation (1) we obtain v_1^1 $2.4v_1^1 - 0.2(0.0845070423) = 1.6$ $2.4v_1^1 = 1.6169014085$ $v_1^1 = 0.6737089202$ Substitute v_1^1 value in equation (3) we obtain v_3^1 $-0.2(0.0845070423) + 2.4v_3^1 = 0.8$

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https://doi.org/10.25130/tjps.v28i4.1532 $2.4v_3^1 = 0.8169014085$ $v_3^1 = 0.3403755869$ Continue to Substitute from 1 = 123, i = 0, m = 0.2in general formula $-0.2v_0^2 + 2.4v_1^2 - 0.2v_2^2$ $= 0.2v_0^1 + 1.6v_1^1 + 0.2v_2^1$ $-0.2v_1^2 + 2.4v_2^2 - 0.2v_3^2$ $= 0.2v_1^1 + 1.6v_2^1 + 0.2v_3^1$ $-0.2v_2^2 + 2.4v_3^2 - 0.2v_4^2$ $= 0.2v_2^1 + 1.6v_3^1 + 0.2v_4^1$ Substitute initial and boundary condition, v_1^1 , v_2^1 , v_3^1 $-0.2(8) + 2.4v_1^2 - 0.2v_2^2$ = 0.2(8) + 1.6(0.6737089202)+ 0.2(0.0845070423) $-0.2v_1^2 + 2.4v_2^2 - 0.2v_3^2$ = 0.2(0.6737089202)+ 1.6(0.0845070423)+0.2(0.3403755869) $-0.2v_2^2 + 2.4v_3^2 - 0.2(4)$ = 0.2(0.0845070423)+ 1.6(0.3403755869) + 0.2(4)Hence, we have $2.4v_1^2 - 0.2v_2^2 = 4.294835681 \dots \dots \dots (1)$ $-0.2v_1^2 + 2.4v_2^2 - 0.2v_3^2$ = 0.338028169.....(2) $-0.2v_2^2 + 2.4v_3^2 = 2.161502347 \dots \dots (3)$ Multiply equation (2) by 12 and solve it with equation (1) we have $-2.4v_1^2 + 28.8v_2^2 - 2.4v_3^2 = 4.056338028$ $2.4v_1^2 - 0.2v_2^2 = 4.294835681$ $28.6v_2^2 - 2.4v_3^2 = 8.351173709 \dots \dots \dots (4)$ Solve equation (4) and adding to equation (3) $28.6v_2^2 - 2.4v_3^2 = 8.351173709$ $-0.2v_2^2 + 2.4v_3^2 = 2.161502347$ $28.4v_2^2 = 10.51267606$ $v_2^2 = 0.37016465$ Substitute v_2^2 value in equation (1) we obtain v_1^2 $2.4v_1^2 - 0.2(0.37016465) = 4.294835681$

 $2.4v_1^2 = 4.368868611$ $v_1^2 = 1.820361921$ Substitute v_2^2 value in equation (3) we obtain v_3^2 $-0.2(0.1850823249) + 2.4v_3^2 = 2.161502347$ $2.4v_3^2 = 2.235535277$ $v_3^2 = 0.931473032$ Continue to Substitute from 1 = 123, i = 2, m = 0.2in general formula $-0.2v_0^3 + 2.4v_1^3 - 0.2v_2^3$ $= 0.2v_0^2 + 1.6v_1^2 + 0.2v_2^2$ $-0.2v_1^3 + 2.4v_2^3 - 0.2v_3^3$ $= 0.2v_1^2 + 1.6v_2^2 + 0.2v_3^2$ $-0.2v_2^3 + 2.4v_3^3 - 0.2v_4^3$ $= 0.2v_2^2 + 1.6v_3^2 + 0.2v_4^2$ Substitute initial and boundary condition, v_1^1 , v_2^1 , v_3^1 $-0.2(8) + 2.4v_1^3 - 0.2v_2^3$ = 0.2(8) + 1.6(1.820361921)+ 0.2(0.37016465) $-0.2v_1^3 + 2.4v_2^3 - 0.2v_3^3$ = 0.2(1.820361921)+ 1.6(0.37016465)+0.2(0.931473032) $-0.2v_2^3 + 2.4v_3^3 - 0.2(4)$ = 0.2(0.37016465)+ 1.6(0.931473032) + 0.2(4)Hence, we have $2.4v_1^3 - 0.2v_2^3 = 6.186612004 \dots \dots \dots (1)$ $-0.2v_1^3 + 2.4v_2^3 - 0.2v_3^3$ $= 1.14263043 \dots \dots (2)$ $-0.2v_2^3 + 2.4v_3^3 = 3.164389781 \dots \dots (3)$ Multiply equation (2) by 12 and solve it with equation (1) we have $-2.4v_1^3 + 28.8v_2^3 - 2.4v_3^3 = 13.71156516$ $2.4v_1^3 - 0.2v_2^3 = 6.186612004$ $28.6v_2^3 - 2.4v_3^3 = 19.89817717 \dots \dots (4)$ Solve equation (4) and adding to equation (3) $28.6v_2^3 - 2.4v_3^3 = 19.89817717$ $-0.2v_2^3 + 2.4v_3^3 = 3.164389781$

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<u>https://doi.org/10.25130/tjps.v28i4.1532</u> $28.4v_2^3 = 23.06256695$

 $v_2^3 = 0.812062217$ Substitute v_2^3 value in equation (1) we obtain v_1^3 $2.4v_1^3 - 0.2(0.812062217) = 6.186612004$ $2.4v_1^3 = 6.349024447$ $v_1^3 = 2.645426853$

Substitute v_2^3 value in equation (3) we obtain v_3^3 -0.2(0.812062217) + 2.4 v_3^3 = 3.164389781 2.4 v_3^3 = 3.326802225 \rightarrow v_3^3 = 1.386167594

Crank-Nicolson	Exact solution
0.67370892	0.98
0.084507042	0.24
0.340375587	0.398
1.820361921	1.98
0.37016465	0.55
0.931473032	0.6
2.645426853	2.398
0.812062217	1.06
1.386167594	1.4

Table (1) Stability of Crank-Nicolson

7- Conclusion

- 1- the Stability of Crank Nicholson is un condition $\forall r^2$.
- 2- the Stability of Explicit Methods is stable if $r^2 \leq \frac{4-(Dt)^2}{4}$.
- 3- The Explicit Method is faster than Crank
 Nicholson method to get the result (50%).
- 4- The Crank Nicholson is more accuracy the Explicit Method .

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